

bond stretching, bending, torsion etc

easiest if just full 3d eq of motion per particle, (no rigid constraint)

- treat all atoms individually

- take care to use small displacement for tightly bound atoms  
⇒ not so efficient.

rigid constraints

- do explicitly

- for large molecules become difficult, MD easier.

eq. of motion

(unlike MD)

Important point: trial moves need not be physical

Cluster moves.

| as long as the reverse has same prob ⇒ detailed balance

Remember Ising model at low T, ↑↑↑↑↑↑ ↓↓↓↓↓↓

- in principle can't all ↑ in all ↓ are connected, but very unlikely sequence of moves due to high energy barrier

Flipping 1 spin at a time will not make this transition,

(common problem in MC & MD)

MC moves can bypass U-barriers: Cluster moves.

- internal interactions in cluster unchanged
- interaction with environment changes,

⇒ best clusters:  $\Delta E \approx 0 \Rightarrow$  high acceptance.

→ tune acceptance probs

example: Ising model.

$$\begin{array}{c} \uparrow \downarrow \uparrow \downarrow \uparrow \downarrow \\ \uparrow \downarrow \uparrow \downarrow \uparrow \downarrow \end{array} \text{etc. } H = -\sum_{nn} J S_i S_j$$

assume positive.

$$U = (N_a - N_p) J$$

↓ anti-parallel # parallel.

cluster: connecting neighbours

- not if antiparallel,
- prob p if parallel

suppose we flip this subset of cluster, which will send  $N_p' = N_p + \Delta$

$$U' = U - 2J\Delta$$

prob of forward move generated:

backward move

connected spins

$\frac{N_p - N_p'}{N_p}$  broken spins

$$N_a' = N_a - \Delta$$

$$\neq p^{N_p} (1-p)^{N_p - N_p'}$$

violate detailed balance?

$$n'_c = n_c$$

must flip some cluster back

$$P(\text{forward})/P(\text{backward}) = (1-p)^{-\Delta}$$

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## detailed balance

$P(\text{original state}) \xrightarrow{\text{generate}} P(\text{forward})$   $P(\text{accept forward})$   
 $\xleftarrow{\text{suppose } = 1} (3)$   
 $P(\text{new state}) \xrightarrow{\text{generate}} P(\text{backward})$   $P(\text{accept backward})$   
 $\xleftarrow{(3 \text{ not exp}(-\beta U)) \text{ can't be, because } P(\text{forward}) \neq P(\text{backward})}$   
 $\exp(-\beta U) \xrightarrow{\text{exp}(-\beta(U - \epsilon j A))}$

$$P(\text{forward})/P(\text{backward}) = \exp(+\beta_2 \Delta) = (1-p)^{-\Delta}$$

$$\Rightarrow \exp(-\beta z_j) = 1-p$$

$\Rightarrow$  always accept and you speed up convergence!

In general not always possible

Early rejection: figure out early on that a move is doomed and never calculate all of it.

- hard-core interactions
  - strongly repulsive component.  
 $\text{If } \Delta V > \text{Something}$

Ensembles in MC: same as thermostats in MD; get right dist<sup>1200</sup>

- So far, done NVT

## Microcanonical MC

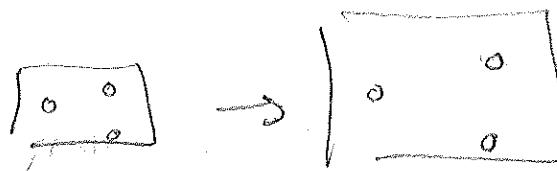
- generate trial move as usual
  - extra variable:  $E_D$
  - modify acceptance
    - if  $\Delta U < 0$  : accept,  $E_D += \Delta U$
    - if  $\Delta U > 0$ ,  $E_D > \Delta U$  : accept  $E_D -= \Delta U$
    - ,  $E_D < \Delta U$  : reject

acceptance not random, but generation of trial move is  
 wiggle around conserved energy.  $U + E_p$  conserved.  
 In principle should include kinetic term as well  
not used much.

## MC in NPT ensemble

- common experimental situation
- construction similar to barostat

rescale the system



$V$  as variable

distribution

$$\int d\mathbf{r} \exp[-\beta(U+PV)] \underset{\substack{\text{Jacobian} \\ \text{scaled coordinate}}}{=} \int d\mathbf{s} V_{\text{new}}^N \exp[-\beta(U+PV)]$$

- ( ) step① trial move,  
either ① normal NVT move  
or ② change in volume:  
 $r \rightarrow r \left( \frac{V_{\text{new}}}{V_{\text{old}}} \right)^{1/3}$

- ( ) step② accept ① according to Metropolis  
or ② according to

$$P = \min \left( 1, \underbrace{\exp[-\beta(\Delta U + P\Delta V)] \left( \frac{V_{\text{new}}}{V_{\text{old}}} \right)^N}_{\tilde{P}} \right)$$

detailed balance

$$P(\rightarrow) \cdot V_{\text{old}}^N \exp[-\beta(U_{\text{old}} + PV)] = P(\leftarrow) \underbrace{V_{\text{new}}^N \exp[-\beta(U_{\text{new}} + PV)]}_{= \min(1, \frac{1}{\tilde{P}})} \Rightarrow \text{true!}$$

end lecture 10